



# MHD Hybrid Nanofluid Flow with Heat Radiation over a Stretching Surface: Numerical Approach

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## Abstract

This study investigates the three-dimensional magnetohydrodynamic (MHD) radiative flow and heat transfer of a kerosene-based hybrid nanofluid containing  $Al_2O_3$  and Cu nanoparticles over a stretching sheet. The governing partial differential equations (PDEs) are formulated and transformed into a coupled system of ordinary differential equations (ODEs) via suitable similarity transformations. The resulting nonlinear ODEs are solved using the Homotopy Analysis Method (HAM), with the effects of various physical parameters on velocity and temperature profiles illustrated through graphical and numerical results. Furthermore, the influences of key parameters—including skin friction coefficients, heat transfer characteristics, rotation parameter, Biot number, and magnetic field strength—on the flow and thermal behaviors are analyzed. The findings reveal that the rotation parameter, Biot number, and magnetic field strength significantly affect the velocity profiles and heat transfer performance of the hybrid nanofluid system.

**Keywords:** hybrid nanofluid, magnetohydrodynamics (MHD), heat radiation, stretching sheet, homotopy analysis method (HAM).

## Special Symbols

Symbol	Description
$\rho$	Density
$\mu$	Dynamic Viscosity
$\nu$	Kinematic Viscosity
$\xi$	Transformed Coordinates
$\phi$	Nanoparticle volume fraction
$\Omega$	Transformed Angular Velocity
$\sigma$	Electrical Conductivity
$\omega$	Angular Velocity
$\gamma$	Biot number

## Subscripts

Symbol	Description
$f$	Fluid
$bf$	Base fluid
$nf$	Nanofluid
$hnf$	Hybrid Nanofluid



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## List of Symbols

Symbol	Description
$T$	Temperature
$u, v, w$	Velocity components
$h$	Heat transfer coefficient
$f, g$	Dimensionless velocities
$\theta$	Dimensionless temperature
$s$	Shape factor
$Nu$	Nusselt number
$k$	Thermal Conductivity
$Re$	Reynolds number
$C_p$	Specific Heat
$R_d$	Radiation parameter
$\tau_{xy}$	Shear stress

## 1 Introduction

A fluid with nanoparticles, particles with sizes in the nanometer range dispersed throughout a base fluid, is called a nanofluid. As the nanoparticles are evenly dispersed throughout the fluid medium, these fluids are essentially colloidal suspensions [1, 2]. Usually found in nanofluids, nanoparticles are composed of metals, metal oxides, carbides, or carbon nanotubes. Water, ethylene glycol, and different types of oil are frequently used as base fluids for these solutions [3].

Nanofluids have a wide range of uses in heat transfer processes, including fuel cells, microelectronics, engine cooling systems, vehicle thermal management, home refrigeration, chillers, heat exchangers, and boiler flue gas grinding, machining, and temperature reduction [4, 5]. Nanofluids have improved convective heat transfer and thermal conductivity when compared to their respective base fluids [6]. Understanding the rheological characteristics of nanofluids is essential for determining their suitability for convective heat transfer applications [7, 8].

When exposed to ultrasonic fields, nanofluids exhibit shear-wave reconversion from incident compressional waves in addition to their unique acoustic characteristics; this effect increases with increasing concentrations of nanoparticles [9]. In simulations of computational fluid dynamics (CFD), nanofluids are frequently handled as single-phase fluids [10, 11]. Nevertheless, a two-phase assumption is used in practically all scholarly works. Classical single-phase fluid theory states that the concentrations and features of a nanofluid's constituent parts determine its physical characteristics [12]. A different

approach is to use a two-phase or two-component framework to model nanofluids [13].

By promoting the dispersion of a nanofluid droplet and creating a structural disjoining pressure close to the contact line, the nanoparticles use diffusion to build a solid-like ordered structure along the contact line [14]. This improvement is not seen, though, for droplets with diameters in the nanometer range since the wetting time scale is substantially shorter than the diffusion time scale [15].

Heat transfer rate is a key factor in many industrial processes' qualitative output. Water and ethylene glycol are examples of common fluids with a restricted potential for heat transmission. To improve the basic fluids thermal characteristics, researchers began incorporating nanoparticles into them. Thus, nanofluids are defined as the suspension of unchanging, uniformly dispersed nanoparticles in base fluids, such as methanol, ethylene glycol, and water ( $H_2O$ ). Each nanofluid has unique properties that depend on the base fluid and these microscopic particles. Metals, metal oxides such as Cu, Ag,  $SiO_2$ ,  $TiO_2$ ,  $Fe_2O_3$ , and carbon nanotubes (MWCNTs and SWCNTs) are among the materials that make up these nanoparticles. The nanoparticles floating in these fluids have an average diameter that is with an average diameter ranging from 1 to 100 nanometers (nm).

The limited thermal conductivity of common base fluids, such as water and ethylene, limits their use in many real-world applications. A new class of fluids known as nanofluids was created to get around this restriction by adding different kinds of nanoparticles to improve their thermal characteristics [16]. Nanofluids are extensively utilized in a wide range of industrial and engineering applications, including solar collectors, microchannel heat sinks, refrigeration systems, electronic device cooling, HVAC systems, and even the food sector, because of their exceptional heat transfer capabilities [17–19]. Research on using nanofluids to improve heat transfer has rapidly expanded during the last ten years.

Although conventional fluids like water and air are commonly used as cooling media, the rate of heat exchange that these fluids can achieve is unsuitable for some sheet materials. For this reason, the use of magnetic fields seems to be acceptable due to its simplicity and non-intrusiveness. Solar collectors are devices that efficiently and sustainably transform solar energy into thermal energy. To increase their capacity to transform the necessary form of energy, these solar

collectors employ a variety of nanofluid types. The use of nanofluids to improve the thermal performance of solar collectors has been the subject of numerous experimental and computational investigations.

## 2 Related Work

By suspending a variety of distinct nanoparticles in a base fluid, researchers created a novel type of nanofluid. Choi et al. [20] were the first to experimentally introduce the idea of nanofluids. In 1995, they found that, in contrast to traditional base fluids, nanofluids improved thermal efficiency. Subsequently, Kang et al. [21] verified these findings experimentally. Afterwards, Lee et al. [22] and Eastman et al. [23] evaluated the thermal conductivity of a water-based nanofluid containing aluminum oxide ( $\text{Al}_2\text{O}_3$ ) and copper (Cu) nanoparticles. Nanofluids' heat transfer performance is greatly improved by their obvious increase in thermal conductivity when compared to conventional base fluids.

Based on the idea of their thermal conductivity, a variety of flow models, as well as the procedures and uses of hybrid nanofluids, have been studied [24]. Between two parallel disks, Ahmed et al. [25] examined the behavior of a nanofluid with  $\text{Al}_2\text{O}_3$  nanoparticles in terms of squeezing flow. Furthermore, heat transmission and skin friction numerical and analytical results were emphasized. In their experimental study of the variation of the heat transfer rate for nanofluid ( $\text{Fe}_3\text{O}_4/\text{H}_2\text{O}$ ) inside horizontal circular tubes under the effect of magnetic fields, Sun et al. [26] found that the intensity of the magnetic field and the heat transfer rate were directly correlated.

Kumar et al. [27] found that the use of nanofluids improved the reliability of electronic chips by 70% and reduced temperature, thermal resistance, and power consumption in a numerical comparison of heat transfer performance between conventional fluids and  $\text{Al}_2\text{O}_3/\text{H}_2\text{O}$  nanofluid systems. Under the influence of an applied magnetic field, Lahmar et al. [28] examined the thermal conductivity and heat transfer rate of ( $\text{Fe}_3\text{O}_4/\text{H}_2\text{O}$ ) nanofluids trapped between two parallel plates. Sheikholeslami and Rokni [29] numerically studied the influence of Coulomb force on nanofluid heat transfer within a porous enclosure, considering thermal radiation effects. The magnetohydrodynamic (MHD) flow of a Casson nanofluid across a vertically positioned plate with convective heating and velocity slip effects was covered by Gbadeyan et al. [30]. Additionally, both numerical and graphical analyses

are performed to examine the impact of radiation phenomena and heat conductivity on the flow.

A new category of nanofluids known as hybrid nanofluids is produced by mixing two or more different kinds of nanoparticles, either in a mixed or compound form. This method's main goal is to maximize the physical and chemical characteristics of various materials in order to create a special fluid with exceptional qualities. These artificial hybrid materials have remarkable qualities that none of the constituent parts could provide on their own. The creation, categorization, and use of diverse nanofluids have been the focus of a substantial amount of research. Performance evaluations of hybrid nanofluids, a new and inventive class of fluid, are now underway. The performance of nanofluids and hybrid nanofluids has been compared in a small number of recent studies [31–34].

An advanced kind of nanofluid called a hybrid or composite nanofluid is created when two or more metals, metal oxides, or a combination of the two are dissolved in a base fluid. The study of how magnetic fields interact with conducting fluids is known as magnetohydrodynamics, or MHD. When a fluid is moving and electrically conductive, currents are induced in reaction to magnetic fields, causing this phenomenon. The magnetic field itself is altered by the forces created by these induced currents acting on the fluid. The Navier-Stokes equations, which describe fluid dynamics, and Maxwell's equations, which explain the behavior of electromagnetic fields, combine to form the governing equations for MHD. The complete MHD phenomenon is captured by a collection of connected differential equations that are formed by these.

Kashi'ie et al. [35] investigated the fluidic system dynamics and heat transfer processes, taking into account the impacts of Joule heating, by computationally analyzing the MHD flow of a hybrid nanofluid ( $\text{Al}_2\text{O}_3/\text{H}_2\text{O}$ ) due to a stretched sheet. Wole-Osho et al. [36] identified the hybrid nanofluid's ( $\text{Al}_2\text{O}_3\text{-Zn}/\text{H}_2\text{O}$ ) flow characteristics and found that the concentration of nanoparticles had a major impact on the flow's viscosity and specific heat. Aly et al. [37] used both theoretical and computational methods to study the MHD stagnation point flow across a stretching sheet of hybrid nanofluid while taking slip and dissipation effects into account. They discovered a relationship between the heat transfer rate and MHD.

After analyzing the rheological performance of

tungsten oxide–engine oil nanofluid at different concentrations and temperatures, Aghahadi et al. [38] found a linear relationship between applied shear stress and shear rate. The Homotopy Analysis Method (HAM) was used by Hayat et al. [39] to study the MHD flow of a hybrid nanofluid ( $\text{Al}_2\text{O}_3\text{-Cu}/\text{H}_2\text{O}$ ) across a revolving stretching sheet. Their research quantitatively examined the effects of different physical parameters on the system's temperature and velocity fields. The results offer important new information about the flow and temperature characteristics of hybrid nanofluids in rotating frames. Huminic et al. [40, 41] examined the rate of heat transfer and entropy creation for both traditional and hybrid nanofluids in a variety of physical settings.

A numerical study of the heat transfer behavior of a hybrid nanofluid passing through an asymmetric channel with porous walls was also carried out by Saba et al. [42]. Additionally, their work used graphical representations to show a variety of important physical impacts. To create hybrid nanoparticles (Di–Ag), de Oliveira et al. [43] carried out an experimental investigation on a unique technique for coating diamond nanoparticles with silver. In a second work, Iqbal et al. [44] examined how the Hall current affected the MHD flow of a hybrid nanofluid in a rotating channel while taking heat radiation and various nanoparticle morphologies into account. The mechanics of heat transmission in nanofluid flows has recently been the subject of numerous investigations [45–53].

Three elements make up the phrase magnetohydrodynamics: dynamics (signaling motion), hydro (signaling fluid, especially water), and magneto (signaling magnetic fields). Hannes Alfvén, who won the 1970 Nobel Prize in Physics for his groundbreaking contributions to the subject, set the groundwork for magnetohydrodynamics (MHD) [54]. These waves were first described by Alfvén as "electromagnetic–hydrodynamic waves." He did, however, propose a more straightforward nomenclature in a subsequent paper, writing: "It may be convenient to call this phenomenon 'magneto-hydrodynamic' waves, as the term 'electromagnetic–hydrodynamic waves' is somewhat complicated" [55].

The core, which is located beneath the Earth's mantle, is made up of a liquid outer core and a solid inner core that are both abundant in iron [56, 57]. The Coriolis

effect causes swirling motions, or eddies, to emerge as the liquid outer core flows in the presence of Earth's magnetic field [58]. The geomagnetic dynamo, a self-sustaining process in which fluid motion produces magnetic fields that strengthen the planet's initial magnetic field, is aided by these eddies [59].

The electromagnetic radiation that matter emits due to the thermal motion of its constituent particles is referred to as thermal radiation. Thermal radiation is released by any material that has a temperature higher than absolute zero. Numerous internal processes, such as electrical excitations, molecular vibrations, and lattice oscillations within the material, are the source of the energy emitted [60]. Charge-acceleration or dipole oscillation transforms kinetic energy into electromagnetism. The majority of thermal radiation at ambient temperature is released in the infrared (IR) spectrum [61]. However, a significant amount of the radiation moves into the visible spectrum as the temperature rises above around  $525\text{ }^\circ\text{C}$  ( $977\text{ }^\circ\text{F}$ ), which results in the material emitting visible light. This phenomenon is called incandescence.

One of the basic processes of heat transport, in addition to conduction and convection, is thermal radiation. Thermal radiation is the main way that heat is transferred from the Sun to the Earth. The sky appears to be blue because of the partial absorption and dispersion of this energy in the atmosphere [62]. The atmosphere carries a large portion of the Sun's radiation to the surface, where it is either reflected or absorbed. Things or occurrences that are normally invisible to the human eye can be detected thanks to thermal radiation. This idea is used by thermographic cameras, which use infrared radiation to create thermal images that show temperature differences between surfaces. These pictures can show a scene's temperature gradient and are frequently used to find objects that are hotter than their surroundings. Animals and humans can be located using infrared imaging because of their body temperature in a dark setting with low levels of visible light. Another type of thermal radiation is cosmic microwave background radiation.

Thermal radiation in idealized systems is studied using the blackbody radiation concept. When a radiating object displays the properties of a perfect black body, this theoretical paradigm is applicable in thermodynamic equilibrium [63]. Planck's law links the body's temperature to its radiative energy production and characterizes the spectral distribution

of blackbody radiation. The Wien's displacement law establishes the wavelength (or frequency) at which a black body's radiation is most intense, whereas the Stefan-Boltzmann law controls the total radiant intensity released by the body [64]. In situations when blackbody radiation is not a reliable approximation, quantum electrodynamics (QED) can be used to represent emission and absorption [65].

One type of electromagnetic radiation that a material emits due to its temperature is called thermal radiation. A semi-analytical method for resolving nonlinear ordinary and partial differential equations is the Homotopy Analysis Method (HAM). It uses the topological idea of homotopy to produce a convergent series solution for nonlinear systems. This is accomplished by employing a homotopy-Maclaurin series to handle the system's nonlinearities. In his PhD dissertation [66], Liao Shijun of Shanghai Jiaotong University originally proposed the HAM in 1992. It was improved upon in 1997 to create a broader form of the homotopy for differential systems. As part of this adjustment, a non-zero auxiliary parameter called the convergence-control parameter was added [67, 68].

### 3 Problem statement

The three-dimensional flow of a hybrid nanofluid across a rotating and stretching surface is examined in this work. The surface stretches in the  $xy$ -plane, and the nanofluid travels in the positive  $z$ -direction ( $z > 0$ ).  $u, v$ , and  $w$  stand for the velocity components in the  $x, y$ , and  $z$  directions, respectively.  $T_f$  and  $T$  denote the fluid's temperature and the surface temperature, respectively. The fluid-surface thermal interaction is described by a heat transfer coefficient,  $h_f$ . Furthermore, as shown in Figure 1, a uniform magnetic field  $B_0$  is applied along the  $z$ -axis.

The following is an expression of the governing equations for energy balance, momentum, and mass conservation:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \quad (1)$$

$$\begin{aligned} & \left[ u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} - 2\omega w \right] \\ & = \left( 1 + \frac{1}{\lambda} \right) \frac{\mu_{\text{hnf}}}{\rho_{\text{hnf}}} \left[ \frac{\partial^2 u}{\partial z^2} - \sigma_{\text{hnf}} B_0^2 u \right], \end{aligned} \quad (2)$$

$$\begin{aligned} & \left[ u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} - 2\omega u \right] \\ & = \left( 1 + \frac{1}{\lambda} \right) \frac{\mu_{\text{hnf}}}{\rho_{\text{hnf}}} \left[ \frac{\partial^2 v}{\partial z^2} - \sigma_{\text{hnf}} B_0^2 v \right], \end{aligned} \quad (3)$$

$$\begin{aligned} & (\rho C_p)_{\text{hnf}} \left[ u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} + w \frac{\partial T}{\partial z} \right] \\ & = \left( K_{\text{hnf}} + \frac{16}{3} \frac{\sigma^* T_\infty^3}{k^*} \right) \frac{\partial^2 T}{\partial z^2} \\ & + \mu_{\text{hnf}} \left( 1 + \frac{1}{\lambda} \right) \left[ \left( \frac{\partial u}{\partial z} \right)^2 + \left( \frac{\partial v}{\partial z} \right)^2 \right], \end{aligned} \quad (4)$$

The boundary conditions are:

$$u = u_w = bx, \quad v = 0, \quad (5)$$

$$-K_{\text{hnf}} \left( \frac{\partial T}{\partial z} \right) = h_f (T_f - T) \quad \text{at } z = 0,$$

$$u \rightarrow 0, \quad v \rightarrow 0, \quad T \rightarrow T_\infty \quad \text{as } z \rightarrow \infty. \quad (6)$$

The hybrid nanofluid in question is made up of copper (Cu) and aluminum oxide ( $\text{Al}_2\text{O}_3$ ) nanoparticles mixed and dissolved in kerosene.  $\phi_1$  and  $\phi_2$  represent the volume fractions of  $\text{Al}_2\text{O}_3$  and Cu nanoparticles, respectively. The sum of  $\phi_1 + \phi_2$  yields the overall concentration of nanoparticles in the hybrid nanofluid,  $\phi_{\text{hnf}}$ . Table 1 lists the thermophysical characteristics of nanoparticles, including density, specific heat, and thermal conductivity. Table 2 lists the thermophysical properties of hybrid nanofluid, including viscosity, specific heat capacity, density, thermal conductivity, and electrical conductivity.  $\rho_f$ ,  $\rho_{s1}$ , and  $\rho_{s2}$ , respectively, denote the densities of the base fluid,  $\text{Al}_2\text{O}_3$ , and Cu nanoparticles.  $(C_p)_f$ ,  $(C_p)_{\text{hnf}}$ ,  $(C_p)_{s1}$  and  $(C_p)_{s2}$  represent the specific heat capacities of the base fluid, the hybrid nanofluid, and the  $\text{Al}_2\text{O}_3$  and Cu particles, respectively.

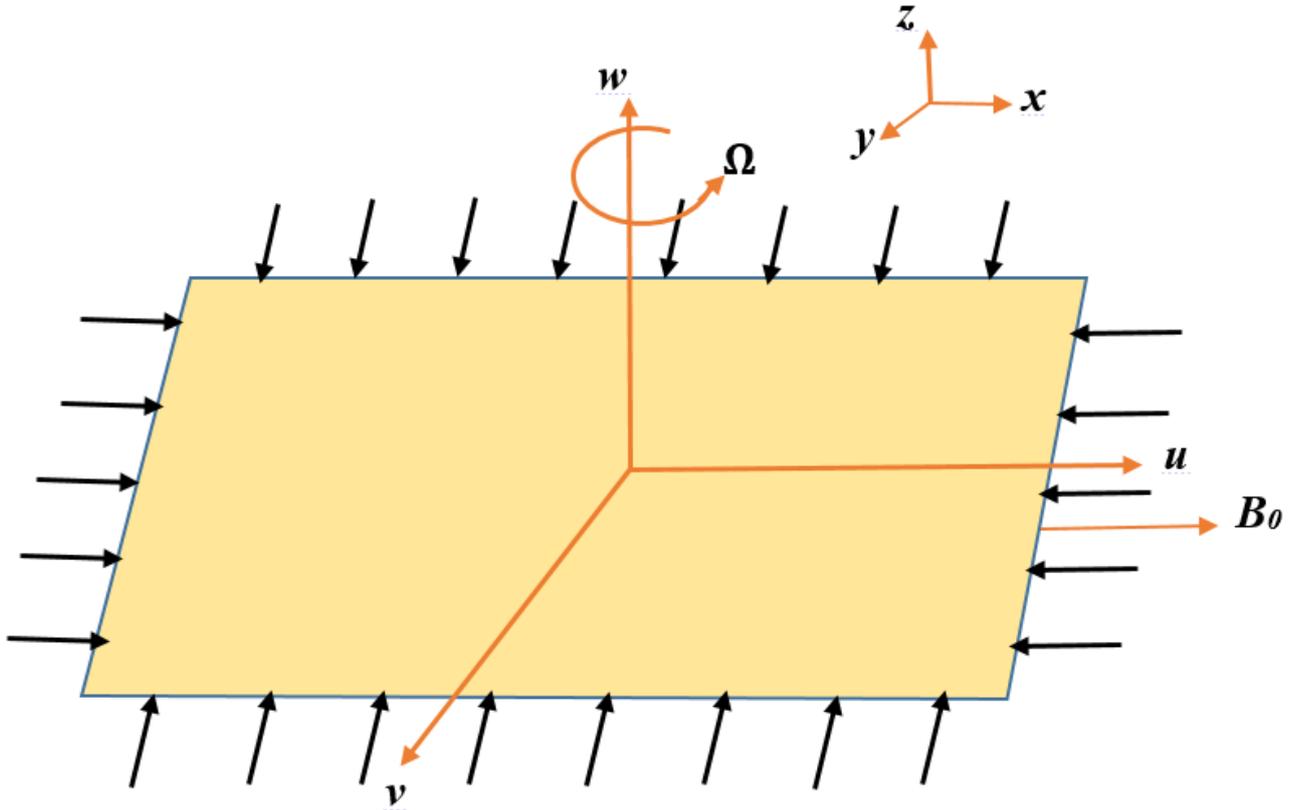
The following are the dimensional expressions for the Nusselt number and the skin friction coefficients:

$$\begin{aligned} C_f &= \frac{\mu_{\text{hnf}}}{\rho_f u_w^2} \left( \frac{\partial u}{\partial z} \right)_{z=0}, \\ C_g &= \frac{\mu_{\text{hnf}}}{\rho_f u_w^2} \left( \frac{\partial v}{\partial z} \right)_{z=0}, \\ Nu &= -\frac{x K_{\text{hnf}}}{k_f (T_f - T_\infty)} \left( \frac{\partial T}{\partial z} \right)_{z=0}. \end{aligned} \quad (7)$$

The following similarity transformations are used to reduce the system of partial differential equations

**Table 1.** Mathematical values of various properties for nanoparticles and base fluid.

Properties	Base fluid		Nano particles	
	Kerosene	Cu	Al <sub>2</sub> O <sub>3</sub>	
Density ( $\rho$ ) (kg m <sup>-3</sup> )	783	8933	3970	
Thermal conductivity ( $K$ ) (W m <sup>-1</sup> K <sup>-1</sup> )	0.145	400	40	
Specific heat ( $C_p$ ) (J kg <sup>-1</sup> K <sup>-1</sup> )	2090	385	765	



**Figure 1.** Geometrical representation of the problem.

(1-4) to a dimensionless set of ordinary differential equations:

$$\begin{aligned}
 u &= bx f'(\xi), & v &= bxg(\xi), & w &= -\sqrt{b\nu_f} f(\xi), \\
 \xi &= \sqrt{\frac{b}{\nu_f}} z, & \theta(\xi) &= \frac{T - T_\infty}{T_f - T_\infty}.
 \end{aligned}
 \tag{8}$$

$$\left(1 + \frac{1}{\beta}\right) g'' - \frac{A_2}{A_1} [f'g - fg' - 2\Omega f' + A_3 M g] = 0,
 \tag{10}$$

$$\begin{aligned}
 [A_4 + R_d] \theta'' + A_5 Pr f \theta' \\
 + \left(1 + \frac{1}{\beta}\right) A_6 Ec (f''^2 + g'^2) = 0
 \end{aligned}
 \tag{11}$$

The system of equations (1-4) is changed into the following form by utilizing the previously established similarity transformations, which also automatically satisfy the continuity equation:

where,

$$\left(1 + \frac{1}{\beta}\right) f''' - \frac{A_2}{A_1} [f'^2 - f f'' - 2\Omega g + A_3 M f'] = 0,
 \tag{9}$$

$$\begin{aligned}
 A_1 &= \frac{\mu_{hnf}}{\mu_f}, & A_2 &= \frac{\rho_{hnf}}{\rho_f}, & A_3 &= \frac{\sigma_{hnf}}{\sigma_f}, \\
 A_4 &= \frac{k_{hnf}}{k_f}, & A_5 &= \frac{(\rho C_p)_{hnf}}{(\rho C_p)_f}
 \end{aligned}
 \tag{12}$$

**Table 2.** Thermophysical properties of hybrid nanofluid.

Properties	Hybrid nanofluid
Viscosity $\mu_{hnf}$	$\frac{\mu_{hnf}}{\mu_f} = \frac{1}{[(1 - \phi_1)(1 - \phi_2)]^{\frac{5}{2}}}$
Specific heat capacity $C_p$	$[\rho C_p]_{hnf} = \left[ (1 - \phi_1) + \frac{(\rho C_p)_{s1}}{(\rho C_p)_f} \phi_1 \right] (1 - \phi_2) + \phi_2 \frac{(\rho C_p)_{s2}}{(\rho C_p)_f}$
Density $\rho_{hnf}$	$\frac{\rho_{hnf}}{\rho_f} = \left[ \left( \frac{\rho_{s1}}{\rho_f} \right) \phi_1 + (1 - \phi_1) \right] (1 - \phi_2) + \phi_2 \frac{\rho_{s2}}{\rho_f}$
Thermal conductivity $k_{hnf}$	$\frac{k_{hnf}}{k_{bf}} = \frac{k_{s2} + k_{bf}(s - 1) - (k_{bf} - k_{s2})(s - 1)\phi_2}{(k_{bf} - k_{s2})\phi_2 + (s - 1)k_{bf} + k_{s2}}$ $\frac{k_{bf}}{k_f} = \frac{k_{s1} + k_f(s - 1) - (k_f - k_{s1})(s - 1)\phi_1}{(k_f - k_{s1})\phi_1 + (s - 1)k_f + k_{s1}}$
Electrical conductivity $\sigma_{hnf}$	$\frac{\sigma_{hnf}}{\sigma_{bf}} = \frac{\sigma_{s2} + 2\sigma_{bf} - 2(\sigma_{bf} - \sigma_{s2})\phi_2}{2\sigma_{bf} + \sigma_{s2} + (\sigma_{bf} - \sigma_{s2})\phi_2}$ $\frac{\sigma_{bf}}{\sigma_f} = \frac{2\sigma_f + \sigma_{s1} - 2(\sigma_f - \sigma_{s1})\phi_1}{2\sigma_f + \sigma_{s1} + (\sigma_f - \sigma_{s1})\phi_1}$

The boundary conditions are:

$$f(\xi) = 0, \quad g(\xi) = 0, \quad f'(\xi) = 1, \quad \theta'(\xi) = -\frac{k_f}{k_{hnf}}\gamma(1 - \theta(\xi)), \quad \text{at } \xi = 0, \quad (13)$$

$$f'(\xi) \rightarrow 0, \quad g(\xi) \rightarrow 0, \quad \theta(\xi) \rightarrow 0, \quad \text{as } \xi \rightarrow \infty. \quad (14)$$

Dimensionless parameters involved in equations (9-11) can be written as:

$$\Omega = \frac{\omega}{b}, \quad M = \frac{\sigma_f B_0^2}{b\rho_f}, \quad Pr = \frac{\mu_f(C_p)_f}{k_f},$$

$$R_d = \frac{16\sigma^* T_\infty^3}{3k^* k_f}, \quad Ec = \frac{b^2 x^2}{(C_p)_f (T_f - T_\infty)}, \quad \gamma = \frac{h_f}{k_f} \sqrt{\frac{\nu_f}{b}} \quad (15)$$

Equation (11), which provides the formulas for the skin friction coefficients and the Nusselt number, can

be expressed as follows:

$$C_f r^{1/2} = \frac{1}{(1 - \phi_1)^{5/2} (1 - \phi_2)^{5/2}} f''(0),$$

$$C_g r^{1/2} = \frac{1}{(1 - \phi_1)^{5/2} (1 - \phi_2)^{5/2}} g'(0), \quad (16)$$

$$Nur^{-1/2} = -\frac{k_{hnf}}{k_f} \theta'(0).$$

where

$$r = \frac{u_w x}{\nu_f}. \quad (17)$$

#### 4 Solution with HAM Method

The Homotopy Analysis Method (HAM), which is incorporated into Mathematica software, is used to numerically solve the transformed system of nonlinear ordinary differential equations, which is shown in Eqs.

(9-11). The final solution is given by the series at  $q = 1$ :

$$\begin{aligned} f(\xi) &= f_0(\xi) + f_1(\xi) + f_2(\xi) + \dots, \\ g(\xi) &= g_0(\xi) + g_1(\xi) + g_2(\xi) + \dots, \\ \theta(\xi) &= \theta_0(\xi) + \theta_1(\xi) + \theta_2(\xi) + \dots \end{aligned} \quad (18)$$

OR

$$\begin{aligned} f(\xi) &= f_0(\xi) + \sum_{m=1}^{\infty} f_m(\xi), \\ g(\xi) &= g_0(\xi) + \sum_{m=1}^{\infty} g_m(\xi), \\ \theta(\xi) &= \theta_0(\xi) + \sum_{m=1}^{\infty} \theta_m(\xi). \end{aligned} \quad (19)$$

$\hbar_f, \hbar_g, \hbar_\theta$  are the auxiliary parameters that control the convergence of the series solutions. These parameters are usually chosen to guarantee both quick and steady convergence of the solution, frequently using graphical analysis employing so-called  $\hbar$ -curves.

### 5 Result and Discussions

In this section, we discussed the effect of various physical parameters on velocity and temperature profile.

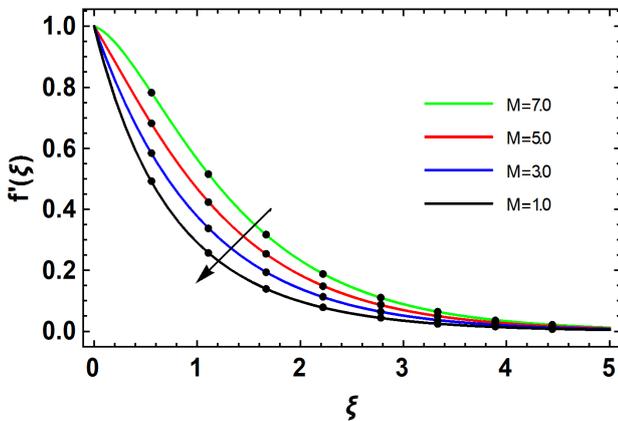


Figure 2. Variation in  $f'(\xi)$  via  $M$ .

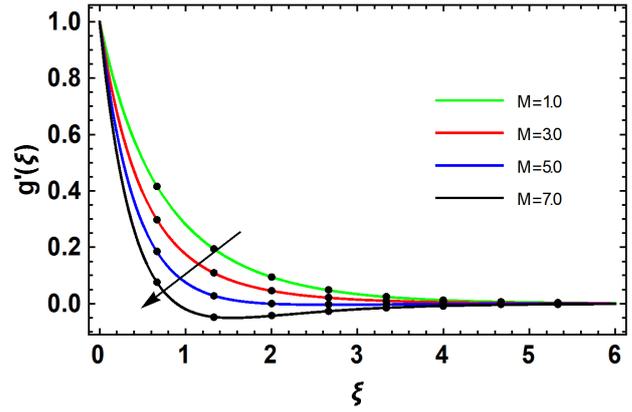


Figure 3. Variation in  $g'(\xi)$  via  $M$ .

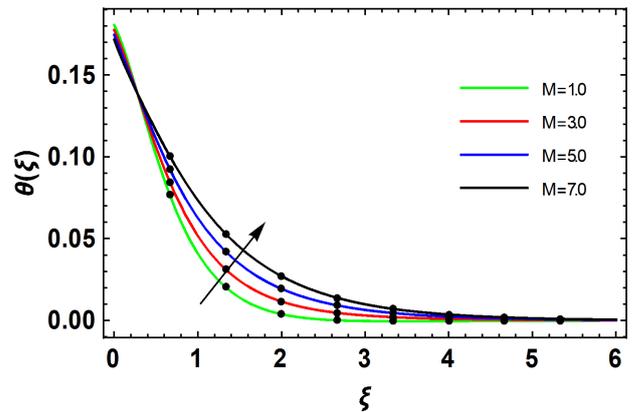


Figure 4. Variation in  $\theta(\xi)$  via  $M$ .

creates resistance. This resistance leads to higher temperatures. Therefore, when  $M$  increases, the temperature profile increases. Figure 4 highlights the effect of  $M$  on temperature profile  $\theta(\xi)$ .

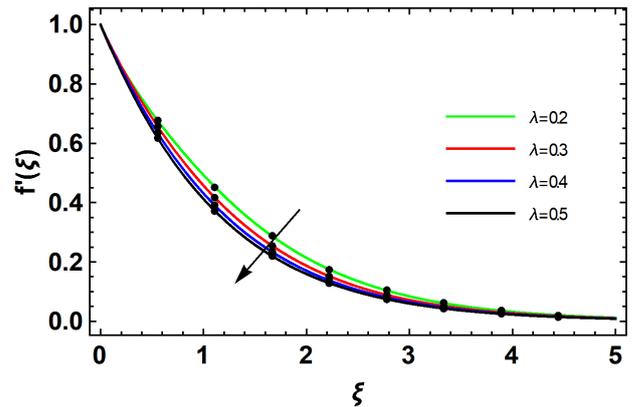


Figure 5. Variation in  $f'(\xi)$  via  $\lambda$ .

Figures 2 and 3 highlights the effect of the magnetic parameter on velocity profiles  $f'(\xi)$  and  $g'(\xi)$ . From figure it is concluded that the velocity profile decreases with higher values of magnetic parameter. Because the increasing values of the magnetic parameter produce a force which is called the Lorentz force, opposite in direction of motion of the fluid. Therefore, when  $M$  increases, the velocity profile decreases. When the magnetic parameter increases, the temperature profile increases. Because magnetic field produces Lorentz force which opposes fluid motion and

Figures 5, 6 and 7 highlights the effect of the Casson parameter  $\beta$  on velocity profiles  $f'(\xi)$ ,  $g'(\xi)$  and temperature profile  $\theta(\xi)$ . The velocity field decreases with an increase in the Casson parameter. This implies that the fluid becomes more viscous with a greater Casson parameter  $\beta$ , which lowers the yield stress and momentum inside the boundary layer. Consequently,

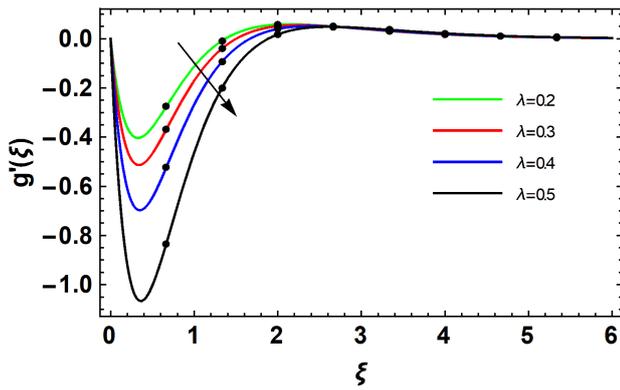


Figure 6. Variation in  $g'(\xi)$  via  $\lambda$ .

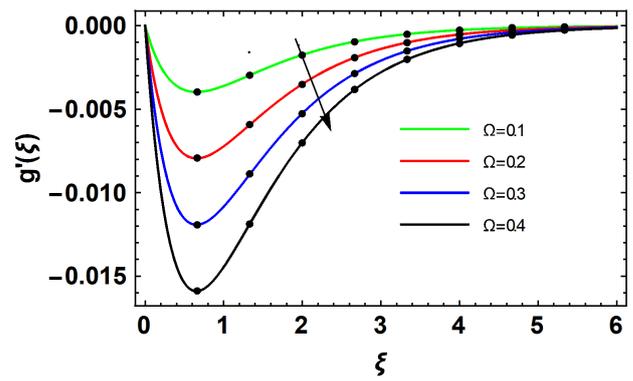


Figure 9. Variation in  $g'(\xi)$  via  $\Omega$ .

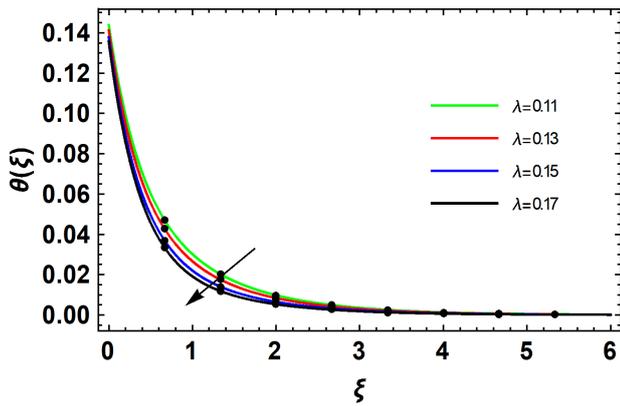


Figure 7. Variation in  $\theta(\xi)$  via  $\lambda$ .

noticeably decrease when  $\Omega$  increases. This happens physically because rotation suppresses the primary and secondary flow velocities by introducing a Coriolis force that acts perpendicular to the flow direction.

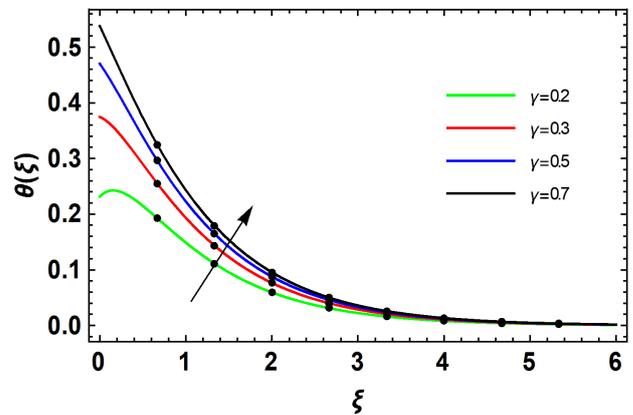


Figure 10. Variation in  $\theta(\xi)$  via  $\gamma$ .

when the value of  $\beta$  rises, the velocity profile decreases as the boundary layer thickness decreases. The way that  $\beta$  behaves on  $\theta(\xi)$  demonstrates that the fluid temperature drops as  $\beta$  rises. This is explained by the fact that a thinner thermal boundary layer is produced when yield stress is reduced due to a greater  $\beta$ . The fluid's temperature drops as the thermal boundary layer gets thinner.

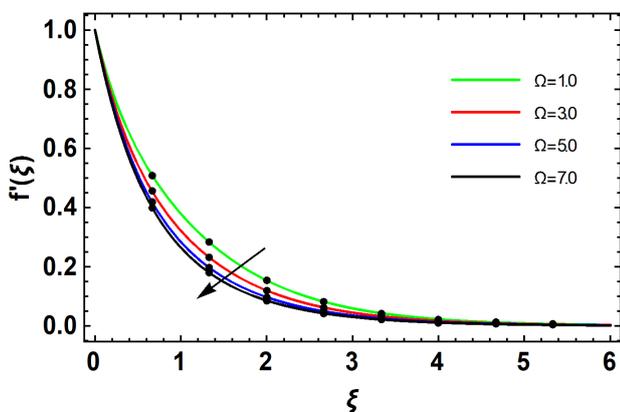


Figure 8. Variation in  $f'(\xi)$  via  $\Omega$ .

For the higher value of parameter  $\gamma$ , the temperature profile  $\theta(\xi)$  increases. Figure 10 highlights the effect of  $\gamma$  on temperature profile  $\theta(\xi)$ . Increasing the thermal Biot number  $\gamma$  usually results in a thicker thermal boundary layer. This is because heat diffusion from the fluid's surface is enhanced by higher convective heat transfer rates, which are associated with larger Biot numbers. As a result, the temperature gradient near the surface decreases, and the thermal boundary layer extends deeper into the fluid.

Figure 11 shows the rise in fluid temperature with the Eckert number  $Ec$ . Frictional heating is responsible for this tendency, in which the fluid's increased viscous dissipation produces more thermal energy. The kinetic energy to the enthalpy difference between the fluid and the boundary surface is represented by the Eckert number  $Ec$ . Due to viscous factors, more kinetic energy is converted into internal energy as  $Ec$  rises, raising the fluid's temperature. As a result, the dimensionless temperature  $\theta(\xi)$  increases with a

The corresponding Figures 8 and 9 demonstrate how the rotational parameter  $\Omega$  affects the velocity profiles  $f'(\xi)$  and  $g'(\xi)$ . Both velocity profiles

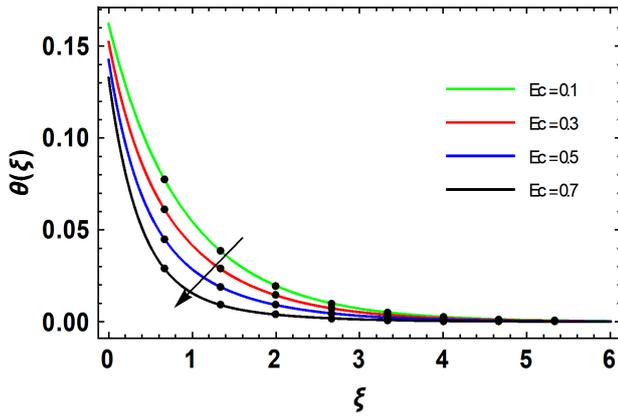


Figure 11. Variation in  $\theta(\xi)$  via  $Ec$ .

higher Eckert number  $Ec$ .

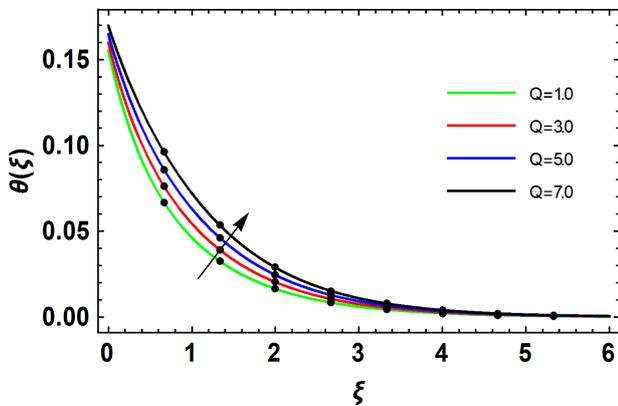


Figure 12. Variation in  $\theta(\xi)$  via  $Q$ .

The impact of the heat source parameter  $Q$  on the temperature distribution  $\theta(\xi)$  is depicted in Figure 12. As  $Q$  rises, the temperature profile shows an upward trend.

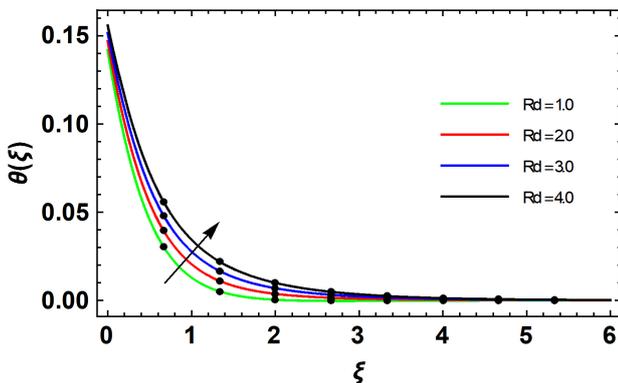


Figure 13. Variation in  $\theta(\xi)$  via  $Rd$ .

The effect of the thermal radiation parameter  $R_d$  on the temperature profile  $\theta(\xi)$  is shown in Figure 13. The graph makes it clear that as  $R_d$  grows, the temperature rises correspondingly. The thickening of the thermal boundary layer, which is caused by stronger thermal

radiation effects at higher values of  $R_d$ , is responsible for this improvement in the thermal profile.

Tables 3, 4 and 5 show the effect of  $\phi_1, \phi_2, \beta, \Omega, M, R_d$  and  $Q$  on skin friction coefficients  $C_f, C_g$  and Nusselt number  $Nu$  for hybrid nanofluid ( $Al_2O_3-Cu/kerosene$ ), respectively.

The computed values of the skin friction coefficient  $C_f r^{1/2}$  for the  $Al_2O_3-Cu/kerosene$  hybrid nanofluid are presented in Table 3. It is observed that as the nanoparticle volume fractions  $\phi_{s1}$  and  $\phi_{s2}$  increase from 0.01 to 0.03, the magnitude of  $C_f$  decreases (from -0.198 to -0.184). An increase in the Casson parameter  $\lambda$  (from 0.6 to 0.9) leads to an increase in the magnitude of  $C_f$ . The rotation parameter  $\Omega$  shows a strong decreasing effect on the magnitude of  $C_f$  as  $\Omega$  increases from 0.2 to 0.5. Similarly, the magnetic parameter  $M$  slightly reduces the magnitude of  $C_f$  with increasing values. In contrast, both the radiation parameter  $R_d$  and heat source parameter  $Q$  cause an increase in the magnitude of  $C_f$  as their values rise.

Table 3. Computed values of the skin friction coefficient  $C_f$  for  $Al_2O_3-Cu/kerosene$  hybrid nanofluid.

$\phi_{s1}$	$\phi_{s2}$	$\lambda$	$\Omega$	$M$	$R_d$	$Q$	$Al_2O_3-Cu/kerosene$
0.01	0.01	0.6	0.2	0.1	0.1	1.0	-0.198246
0.02	0.02						-0.192519
0.03	0.03						-0.184477
		0.7					-0.190192
		0.8					-0.198463
		0.9					-0.202453
			0.3				-0.158890
			0.4				-0.146411
			0.5				-0.135737
				0.3			-0.173476
				0.5			-0.173312
				0.7			-0.173169
					0.2		-0.176343
					0.3		-0.180122
					0.4		-0.183451
						3.0	-0.187345
						5.0	-0.190021
						7.0	-0.195532

Table 4 displays the values of the skin friction coefficient  $C_g r^{1/2}$ . The magnitude of  $C_g$  increases with higher nanoparticle volume fractions  $\phi_{s1}$  and  $\phi_{s2}$  (from -0.487 to -0.539). The Casson parameter  $\lambda$  exhibits a non-monotonic behavior, but overall tends to reduce the magnitude in some ranges. Increasing the rotation parameter  $\Omega$  and magnetic parameter  $M$  both lead to a decrease in the magnitude of  $C_g$ . The heat source parameter  $Q$  shows an increasing effect on the magnitude of  $C_g$  as  $Q$  rises from 1.0 to 7.0.

**Table 4.** Computed values of the skin friction coefficient  $C_g$  for  $Al_2O_3$ -Cu/kerosene hybrid nanofluid.

$\phi_{s1}$	$\phi_{s2}$	$\lambda$	$\Omega$	$M$	$Q$	$Al_2O_3$ -Cu/kerosene
0.01	0.01	0.6	0.2	0.1	1.0	-0.486627
0.02	0.02					-0.511967
0.03	0.03					-0.538907
		0.7				-0.868212
		0.8				-0.567374
		0.9				-0.567066
			0.3			-0.522488
			0.4			-0.484038
			0.5			-0.450855
				0.3		-0.533451
				0.5		-0.510266
				0.7		-0.488345
					3.0	-0.476243
					5.0	-0.512832
					7.0	-0.543451

The local Nusselt number  $N_{ur}^{-1/2}$  values for the  $Al_2O_3$ -Cu/kerosene hybrid nanofluid are shown in Table 5. The results indicate that the Nusselt number increases with increasing nanoparticle volume fractions  $\phi_{s1}$  and  $\phi_{s2}$  (from 0.179 to 0.211), suggesting enhanced heat transfer rate due to the presence of hybrid nanoparticles. Similarly, higher values of the Casson parameter  $\lambda$ , rotation parameter  $\Omega$ , magnetic parameter  $M$ , radiation parameter  $R_d$ , and heat source parameter  $Q$  all contribute to an increase in the Nusselt number, implying improved convective heat transfer performance under these conditions.

**Table 5.** Computed values of the Nusselt number  $Nu$  for  $Al_2O_3$ -Cu/kerosene hybrid nanofluid.

$\phi_{s1}$	$\phi_{s2}$	$\lambda$	$\Omega$	$M$	$R_d$	$Q$	$Al_2O_3$ -Cu/kerosene
0.01	0.01	0.6	0.2	0.1	0.1	1.0	0.179427
0.02	0.02						0.194138
0.03	0.03						0.211145
		0.7					0.190192
		0.8					0.198993
		0.9					0.209453
			0.3				0.158899
			0.4				0.178854
			0.5				0.199847
				0.3			0.173476
				0.5			0.177653
				0.7			0.180021
					0.2		0.176343
					0.3		0.180122
					0.4		0.183651
						3.0	0.182345
						5.0	0.190121
						7.0	0.195512

## 6 Conclusion

In conclusion, the three-dimensional magnetohydrodynamics (MHD) radiative flow of a hybrid nanofluid over a stretching sheet was thoroughly investigated in this study using the homotopy analysis method, which yielded important insights into the intricate interactions between the physical parameters controlling such systems. The work demonstrated the strong influence of important variables, such as rotation parameter, biot number, and magnetic field strength, on the temperature and velocity distributions inside the nanofluid flow by deriving convergent series solutions. The results show that adjusting these factors can dramatically change the thermal properties and flow dynamics, providing a technique to maximize heat transfer in cutting-edge engineering applications. In addition to improving theoretical knowledge of hybrid nanofluid behavior under the influence of radiative and magnetic forces, the methodology used in this work creates a strong analytical foundation for forecasting and managing thermal performance in real-world scenarios. Furthermore, the results indicate that more research into the application of various nanoparticle combinations, different base fluids, and different flow configurations may increase the usefulness and effectiveness of hybrid nanofluids in technological and industrial processes, opening the door for further advancements in energy systems and thermal management.

## Data Availability Statement

Data will be made available on request.

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## Conflicts of Interest

The authors declare no conflicts of interest.

## AI Use Statement

The authors declare that no generative AI was used in the preparation of this manuscript.

## Ethical Approval and Consent to Participate

Not applicable.

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